A Note on the Expansion of the First Order Melnikov Function Near a Class of 3-polycycles^{*}

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Abstract This paper comments that there exist some mistakes in the asymptotic expansion of the first order Menikov function near a 3-polycycle given by Theorem 3.1 of [2]. We present a correction to the theorem, and then use it to show that only one limit cycle can be found near a 3-polycycle for a class of quadratic systems.

Keywords Limit cycle, Quadratic polynomial system, 3-polycycle.

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Consider a system of the form

$$\dot{x} = F(x)y + \epsilon p(x,y), \quad \dot{y} = G(x) + R(x)y^2 + \epsilon q(x,y) \tag{1}$$

with

$$p(x,y) = \sum_{j=0}^{n} \hat{p}_j(x)y^j, \quad q(x,y) = \sum_{j=0}^{n} \hat{q}_j(x)y^j$$

where $\epsilon > 0$ is a small parameter, F(x), G(x), R(x), $\hat{p}_j(x)$ and $\hat{q}_j(x)$ are C^{∞} functions in their variable x.

Regarding system (1), we make two assumptions below as in [2]:

- (A1) F(0) = 0, $F(x) = xF_1(x)$, $F_1(0) > 0$, G(0) > 0, R(0) < 0;
- (A2) System (1) $|_{\epsilon=0}$ has a 3-polycycle L_0 passing through the saddle points $(0, \pm y_0)$ and $(x_1, 0)$, where $y_0 > 0$, $x_1 > 0$. Please see Figure 1.

Then, by Lemma 2.1 of [2], one can see that system $(1)|_{\epsilon=0}$ has an integrating factor of the form

$$\mu(x) = x^{\alpha}\mu_0(x), \ \alpha = -\frac{2R(0)}{F_1(0)} - 1,$$

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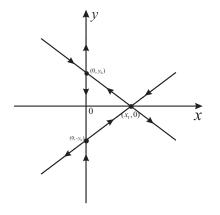


Figure 1. A 3-polycycle.

where $\mu_0(x) \in C^{\infty}$, $\mu_0(0) = \frac{1}{F_1(0)} > 0$. Obviously, under the above assumptions, system $(1)|_{\epsilon=0}$ has a family of periodic orbits inside L_0 given by L_h , $0 < -h < h_0$. Then, correspondingly we have the first order Melnikov function for system (1)

$$M(h) = \oint_{L_h} \mu(x)q(x,y)dx - \mu(x)q(x,y)dy, \quad 0 < -h < h_0.$$
⁽²⁾

Theorem 3.1 of [2] says that the above function has an asymptotic expansion of the form for $0<-h\ll 1$

$$M(h) = \sum_{i \ge 0, \beta_i \notin \mathbb{Z}^+} B_i |h|^{\beta_i} + \left(\sum_{i \ge 0, \beta_i \in \mathbb{Z}^+} B_i |h|^{\beta_i} + \sum_{i \ge 0} C_{2i+1} |h|^{i+1}\right) \ln |h| + \sum_{i \ge 0} b_i h^i.$$
(3)

However, we recently found that the term $C_{2i+1}|h|^{i+1}$ in (3) should be corrected as $C_{2i+1}h^{i+1}$. The reason is that the authors used $|h|^{i+1}$ for h^{i+1} in (3.8) of [2] by mistake. In fact, according to Theorem 3.2.9 of [1], the formula in (3.8) in page 374 of [2] should be corrected as

$$I_2(h) = \sum_{i \ge 0} C_{2i+1} h^{i+1} \ln |h| + N_1(h).$$

Therefore, about the expansion of (2) near h = 0, one has **Theorem 1.** Under (A1) and (A2) with $\alpha \in (-1,0) \cup (0,+\infty)$, we have the asymptotic expansion below

$$M(h) = \sum_{i \ge 0, \beta_i \notin \mathbb{Z}^+} B_i |h|^{\beta_i} + \Big(\sum_{i \ge 0, \beta_i \in \mathbb{Z}^+} B_i |h|^{\beta_i} + \sum_{i \ge 0} C_{2i+1} h^{i+1} \Big) \ln |h| + \sum_{i \ge 0} b_i h^i h^{i+1} \Big) \ln |h| + \sum_{i \ge 0} b_i h^{i+1} h^{i+1} \Big) \ln |h| + \sum_{i \ge 0} b_i h^{i+1} h^{i+1} \Big) \ln |h| + \sum_{i \ge 0} b_i h^{i+1} h^{i+1} \Big) \ln |h| + \sum_{i \ge 0} b_i h^{i+1} h^{i+1} h^{i+1} \Big) \ln |h| + \sum_{i \ge 0} b_i h^{i+1} h^{i+1$$

where $0 < -h \ll 1$, $\beta_i = \frac{\alpha+i}{\alpha+1}$ and B_i , C_i are constants defined in [2]. Then, Corollary 3.2 of [2] should be corrected as the following

Corollary 1. Under (A1) and (A2) with $\alpha \in (-1,0) \cup (0,+\infty)$, then for $0 < -h \ll 1$

(1) if $\alpha \in (-1, 0)$, then

$$M(h) = B_0 |h|^{\frac{\alpha}{\alpha+1}} + b_0 + (B_1 - C_1)|h| \ln|h| + b_1 h + C_3 h^2 \ln|h| + O(h^2);$$

(2) if $\alpha \in (0,1)$, then

$$M(h) = b_0 + B_0 |h|^{\frac{\alpha}{\alpha+1}} + (B_1 - C_1)|h| \ln|h| + b_1 h + B_2 |h|^{\frac{\alpha+2}{\alpha+1}} + C_3 h^2 \ln|h| + O(h^2);$$

(3) if $\alpha \in [1, +\infty) \setminus \{m-2, m \ge 3, m \in \mathbb{N}^+\}$, then

$$M(h) = b_0 + B_0 |h|^{\frac{\alpha}{\alpha+1}} + (B_1 - C_1)|h|\ln|h| + b_1 h + \sum_{i=2}^{[\alpha+2]} B_i |h|^{\frac{\alpha+i}{\alpha+1}} + C_3 h^2 \ln|h| + O(h^2);$$

(4) if $\alpha = m - 2, m \ge 3, m \in \mathbb{N}^+$, then

$$M(h) = b_0 + B_0 |h|^{\frac{\alpha}{\alpha+1}} + (B_1 - C_1)|h|\ln|h| + b_1 h + \sum_{i=2}^{m-1} B_i |h|^{\frac{\alpha+i}{\alpha+1}} + (B_m + C_3)|h|^2 \ln|h| + O(h^2).$$

Now, we present the formulas for b_0, b_1, C_1, B_0, B_1 . Note that system (1) is a near-integrable system with the first integral of the form

$$H(x,y) = x^{\alpha+1} \Big[P_1(x) + \frac{1}{2} \mu_0(x) F_1(x) y^2 \Big],$$

where

$$x^{\alpha+1}P_1(x) = -\int \mu(x)G(x)dx \in C^{\infty}.$$

Further, we suppose that the expansions of $\mu_0(x)$, F(x), G(x), $P_1(x)$, p(x,y) and q(x,y) at the origin are of the form respectively

$$\begin{split} \mu_0(x) &= \sum_{j \ge 0} a_j x^j, \quad F(x) = \sum_{j \ge 1} f_j x^j, \quad G(x) = \sum_{j \ge 0} g_j x^j, \quad P_1(x) = \sum_{j \ge 0} p_j x^j, \\ p(x,y) &= \sum_{j=0}^n \sum_{i \ge 0} p_{ij} x^i y^j, \quad q(x,y) = \sum_{j=0}^n \sum_{i \ge 0} q_{ij} x^i y^j. \end{split}$$

Meanwhile, suppose that H(x, y) at the point $(x_1, 0)$ can be expanded as

$$H(x,y) = \frac{\lambda}{2} \left[y^2 - (x-x_1)^2 \right] + \sum_{i \ge 1} h_{i2} (x-x_1)^i y^2 + \sum_{i \ge 3} h_{i0} (x-x_1)^i, \ \lambda = |\mu_1(x_1) \sqrt{F(x_1)G'(x_1)}|.$$

Clearly, $\mu(x)p(x,y)$ and $\mu(x)q(x,y)$ can be expressed as at the point $(x_1,0)$

$$\mu(x)p(x,y) = \sum_{j=0}^{n} \sum_{i \ge 0} a_{ij}(x-x_1)^i y^j, \quad \mu(x)q(x,y) = \sum_{j=0}^{n} \sum_{i \ge 0} b_{ij}(x-x_1)^i y^j.$$

Then, by [2], we have

$$\begin{split} C_1 &= -\frac{1}{\lambda} \Big[(\mu p)_x + (\mu q)_y \Big] (x_1, 0), \\ C_3 &= \frac{-1}{2\lambda^2} \Big[(-3a_{30} - b_{21} + a_{12} + 3b_{03}) - \frac{1}{\lambda} (2a_{20}) + b_{11}) (3h_{30} - h_{12}) \Big], \\ b_0 &= \begin{cases} \lim_{h \to 0^-} \oint_{L_h} \mu q dx - \mu p dy + O(B_0) & \alpha \in (-1, 0), \\ \lim_{h \to 0^-} \oint_{L_h} \mu q dx - \mu p dy, & \alpha \in (0, +\infty), \end{cases} \end{split}$$

$$\begin{split} b_{1}|_{B_{0}=B_{1}-C_{1}=0} &= \lim_{h \to 0^{-}} \oint_{L_{h}} \left[(\mu p)_{x} + (\mu q)_{y} \right] dt, \\ B_{0} &= \sum_{j=0}^{\left[\frac{n}{2}\right]} \frac{2^{j+\frac{3}{2}}}{2j+1} |p_{0}|^{-\frac{\alpha_{j}+1}{\alpha+1}} \alpha a_{0} p_{0,2j} B_{0j}, \\ B_{1} &= \sum_{j=0}^{\left[\frac{n}{2}\right]} \frac{2^{j+\frac{3}{2}}}{2j+1} |p_{0}|^{-\frac{\alpha_{j}}{\alpha+1}} \left[\frac{\alpha_{j}+2}{\alpha+1} |p_{0}|^{-\frac{\alpha+3}{\alpha+1}} p_{1} - \frac{1}{2} (2j+1)(a_{1}f_{1} + a_{0}f_{2}) |p_{0}|^{-\frac{2}{\alpha+1}} \right] \alpha a_{0} p_{0,2j} B_{1j} \\ &+ \sum_{j=0}^{\left[\frac{n}{2}\right]} \frac{2^{j+\frac{3}{2}}}{2j+1} |p_{0}|^{-\frac{\alpha_{j}+2}{\alpha+1}} \left[a_{0}q_{0,2j+1}(2j+1) + (\alpha+1)(a_{0}p_{1,2j} + a_{1}p_{0,2j}) \right] B_{1j} \end{split}$$

where B_{ij} 's are constants satisfying $B_{0j} > 0$ (resp. < 0) for $\alpha \in (-1,0)$ (resp. $\alpha \in (0, +\infty)$), and $\alpha_j = \frac{\alpha-3}{2} - (\alpha+1)j$.

Further, the authors of [2] using (3) (i.e incorrect results) studied a quadratic system of the from

$$\dot{x} = axy + \epsilon \sum_{i+j=0}^{2} p_{ij} x^{i} y^{j}, \quad \dot{y} = b + cx + dx^{2} + ey^{2} + \epsilon \sum_{i+j=0}^{2} q_{ij} x^{i} y^{j}, \qquad (4)$$

where $\epsilon > 0$ sufficiently small, a, b, d > 0, c, e < 0, and $\frac{b}{e} = -\frac{c^2(a-e)}{d(a-2e)^2}$. The unperturbed system (4) $|_{\epsilon=0}$ has a 3-polycycle. From Theorem 5.1 of [2], the authors [2] claimed to find at least 2 limit cycles near the 3-polycycle for system (4). In fact, by using the corrected result, one can only get one limit cycle. Here, we give the detailed proof.

One can easily see that system $(4)|_{\epsilon=0}$ is an integrable system having an integrating factor $\mu(x) = \frac{1}{a}x^{\alpha}$, and it has a first integral given by

$$H(x,y) = \frac{x^{\alpha+1}}{2} \Big[y^2 - \frac{d}{a-e} \Big(x - \frac{b(a-2e)}{ce} \Big)^2 \Big], \quad \alpha = -1 - \frac{2e}{a}$$

Further it has a center at $\left(\frac{b(a-2e)}{c(e-a)}, 0\right)$ and three elementary saddle points at $\left(0, \pm \sqrt{\frac{-b}{e}}\right)$ and $\left(\frac{b(a-2e)}{ce}, 0\right)$, respectively. There exists a family of periodic orbits given by

$$L_h: H(x,y) = h, h \in (\eta, 0),$$

around the center and bounded by a 3-polycycle through the three saddle points. Here, $\eta = -\frac{1}{2} \frac{d}{a-e} \frac{a^2}{e^2} \left[\frac{b(a-2e)}{c(e-a)} \right]^{\frac{2(a-e)}{a}} < 0$. Corresponding to system (4), we have the following first order Melnikov function

$$M(h) = \oint_{L_h} \frac{1}{a} x^{\alpha} \sum_{i+j=0}^{2} q_{ij} x^i y^j dx - \frac{1}{a} x^{\alpha} \sum_{i+j=0}^{2} p_{ij} x^i y^j dy$$

By using integration by parts, it follows that

$$M(h) = \frac{1}{a} \oint_{L_h} x^{\alpha} \Big\{ q_{01} + (\alpha + 1)p_{10} + \Big[q_{11} + (\alpha + 2)p_{20} \Big] x + \frac{\alpha p_{00}}{x} \Big\} y dx + \frac{\alpha}{3a} p_{02} \oint_{L_h} x^{\alpha - 1} y^3 dx$$
$$= A_0 \oint_{L_h} x^{\alpha - 1} y dx + A_1 \oint_{L_h} x^{\alpha} y dx + A_2 \oint_{L_h} x^{\alpha + 1} y dx, \tag{5}$$

since $y^2 = 2hx^{-\alpha-1} + \frac{d}{a-e} \left[x - \frac{b(a-2e)}{xe} \right]^2$ along the curve L_h . Here,

$$A_{0} = \frac{\alpha}{a}p_{00} + \frac{b\alpha}{a(a-e)}p_{02},$$

$$A_{1} = \frac{1}{a}q_{01} + \frac{\alpha+1}{a}p_{10} + \frac{c\alpha}{a(a-e)}p_{02},$$

$$A_{2} = \frac{1}{a}q_{11} + \frac{\alpha+2}{a}p_{20} + \frac{d\alpha}{a(a-e)}p_{02}.$$

About these integrals in (5), one finds that **Lemma 1.** For $h \in (\eta, 0)$, we have

$$\oint_{L_h} x^{\alpha} y dx = \frac{c(e-a)}{b(a-2e)} \oint_{L_h} x^{\alpha+1} y dx.$$

Proof. Note that the curve L_h can be rewritten as $y = \pm \sqrt{2hx^{-\alpha-1} + \frac{d}{a-e} \left[x - \frac{b(a-2e)}{ce}\right]^2}$ for $x_1(h) \le x \le x_2(h)$, where $x_1(h)$ and $x_2(h)$ are the solutions of the equation $-\frac{d}{a-e} \frac{x^{\alpha+1}}{2} \left[x - \frac{b(a-2e)}{ce}\right]^2 = h$ satisfying $0 < x_1(h) < \frac{b(a-2e)}{c(e-a)} < x_2(h)$. Then,

$$\frac{c(e-a)}{b(a-2e)} \oint_{L_h} x^{\alpha+1} y dx - \oint_{L_h} x^{\alpha} y dx
= \oint_{L_h} x^{\alpha} \Big[\frac{c(e-a)}{b(a-2e)} x - 1 \Big] y dx
= 2 \int_{x_1(h)}^{x_2(h)} x^{\alpha} \Big[\frac{c(e-a)}{b(a-2e)} x - 1 \Big] \sqrt{2hx^{-\alpha-1} + \frac{d}{a-e} [x - \frac{b(a-2e)}{ce}]^2} dx
= 2 \int_{x_1(h)}^{x_2(h)} x^{\frac{\alpha-1}{2}} \Big[\frac{c(e-a)}{b(a-2e)} x - 1 \Big] \sqrt{\frac{d}{a-e}} x^{\alpha+1} \Big[x - \frac{b(a-2e)}{ce} \Big]^2 + 2h dx.$$
(6)

Let $u = x^{\frac{\alpha+1}{2}} \left[x - \frac{b(a-2e)}{ce} \right]$. Then, $du = \frac{\alpha+1}{2} \frac{b(a-2e)}{ce} x^{\frac{\alpha-1}{2}} \left[\frac{c(e-a)}{b(a-2e)} x - 1 \right] dx$, and the function u is monotonic for $x \in (x_1(h), \frac{b(a-2e)}{c(e-a)})$ (resp. $x \in (\frac{b(a-2e)}{c(e-a)}, x_2(h))$). Then, by the above discussion, one can derive that from (6)

$$\begin{aligned} &\frac{c(e-a)}{b(a-2e)} \oint_{L_h} x^{\alpha+1} y dx - \oint_{L_h} x^{\alpha} y dx \\ &= \frac{4}{\alpha+1} \frac{ce}{b(a-2e)} \Big[\int_{\left[\frac{2(e-a)h}{d}\right]^{\frac{1}{2}}}^{\rho} \sqrt{\frac{d}{a-e}u^2 + 2h} dx + \int_{\rho}^{\left[\frac{2(e-a)h}{d}\right]^{\frac{1}{2}}} \sqrt{\frac{d}{a-e}u^2 + 2h} dx \Big] \\ &\equiv 0, \end{aligned}$$

where $\rho = \frac{ab(a-2e)}{ec(e-a)} \left[\frac{b(a-2e)}{c(e-a)}\right]^{\frac{\alpha+1}{2}}$. This ends the proof. Thus, by Lemma 1, the function in (5) can be expressed as

$$M(h) = A_0 \oint_{L_h} x^{\alpha - 1} y dx + \left[A_1 + \frac{b(a - 2e)}{c(e - a)} A_2 \right] \oint_{L_h} x^{\alpha} y dx.$$
(7)

We note that $\alpha > -1$ and $\alpha \neq 0$. By Corollary 1, we should discuss the expansion of M(h) for $\alpha \in (-1,0)$ or $\alpha \in (0,+\infty)$, separately. Here, we only provide a proof

for the case $\alpha \in (0, +\infty)$ since the proof for the case $\alpha \in (-1, 0)$ is similar. For $\alpha > 0$, from Corollary 1 again, the first order Melnikov function M(h) in (7) near the 3-polycycle can be expanded as

$$M(h) = b_0 + B_0 |h|^{\frac{\alpha}{\alpha+1}} + O(|h|\ln|h|), \quad 0 < -h \ll 1.$$
(8)

It is easy to see that

$$B_{0} = 2\sqrt{2} \left(-\frac{b}{2e}\right)^{\frac{1-\alpha}{2(1+\alpha)}} B_{00} \left[\frac{\alpha}{a} p_{00} + \frac{b\alpha}{a(a-e)} p_{02}\right]$$
$$= 2\sqrt{2} \left(-\frac{b}{2e}\right)^{\frac{1-\alpha}{2(1+\alpha)}} B_{00} A_{0}.$$

Then, we have from (8)

$$b_0 = \lim_{h \to 0^-} M(h) = \lim_{h \to 0^-} \oint_{L_h} \mu q dx - \mu p dy = \oint_{L_0} \mu q dx - \mu p dy$$

which, implies that

$$b_{0} = -2A_{0}\sqrt{\frac{d}{a-e}} \int_{0}^{\frac{b(a-2e)}{ce}} x^{\alpha-1} \left[x - \frac{b(a-2e)}{ce}\right] dx - 2\left[A_{1} + \frac{b(a-2e)}{c(e-a)}A_{2}\right] \\ \times \sqrt{\frac{d}{a-e}} \int_{0}^{\frac{b(a-2e)}{ce}} x^{\alpha} \left[x - \frac{b(a-2e)}{ce}\right] dx \\ = \frac{2A_{0}}{\alpha(\alpha+1)}\sqrt{\frac{d}{a-e}} \left[\frac{b(a-2e)}{ce}\right]^{\alpha+1} + \frac{2\left[A_{1} + \frac{b(a-2e)}{c(e-a)}A_{2}\right]}{(\alpha+1)(\alpha+2)}\sqrt{\frac{d}{a-e}} \left[\frac{b(a-2e)}{ce}\right]^{\alpha+2}.$$

Obviously, B_0 and b_0 can be taken as free parameters. We can choose them satisfying $0 < -b_0 \ll B_0 \ll 1$ such that the sign of the function M(h) in (8) near h = 0 can be changed one time. This means that one can find only one limit cycle near the 3-polycycle. Further, as $b_0 = B_0 = 0$, we have

$$A_0 = 0, \quad A_1 + \frac{b(a-2e)}{c(e-a)}A_2 = 0$$

Hence, when $b_0 = B_0 = 0$, we have $M(h) \equiv 0$ for $h \in (\eta, 0)$. Thus, one can find one limit cycle near the 3-polycycle by using the first order Melnikov function.

References

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